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## Crystal Structure

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# $\alpha$-Tris(2,4-pentanedionato- $\kappa^{2} O, O^{\prime}$ )cobalt(III) at 240, 210, 180, 150 and 110 K 

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The crystal structure of the title compound, $\left[\mathrm{Co}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$, has been investigated by a multi-temperature measurement. In contrast to the isomorphous Al compound, the title compound exists in the studied temperature range as its monoclinic $\alpha$ polymorph (space group $P 2_{1} / c$ ) and does not undergo a phase transition. Rigid-body TLS analyses have been performed and the anisotropic thermal expansion tensor $\alpha_{i j}$ has been determined. The cell axes show a linear expansion behavior with respect to the temperature, but the slope is significantly different. A possible explanation are the different strengths of different intermolecular $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ contacts, which run in different crystallographic directions.

## Comment

Metal complexes, especially transition metal complexes, play an important role in catalytic processes. Trivalent metalacetylacetonate (2,4-pentanedionate, acac) complexes are particularly accessible species for studying structure, bonding and ligand coordination in organometallic systems. Polymorphism is well known for these complexes [see e.g. Geremia \& Demitri (2005) for $\left.\mathrm{Mn}(\mathrm{acac})_{3}\right]$. In the course of our ongoing studies on these complexes (von Chrzanowski et al., 2006, 2007), we have now determined the crystal structure of $\mathrm{Co}(\mathrm{acac})_{3}$, (I). According to an old nomenclature of Astbury \& Morgan (1926), the $\alpha$ polymorph crystallizes in the monoclinic crystal system with space group $P 2_{1} / c$. In our tempera-ture-dependent study of $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$, which is isomorphous with $\alpha-\mathrm{Co}(\mathrm{acac})_{3}$, we observed a phase transition to a new $\delta$ polymorph. This phase transition occurs between 150 and 110 K .

To our knowledge, all previous crystallographic investigations of $\mathrm{Co}(\mathrm{acac})_{3}$ were carried out at room temperature. Only the $\alpha$ polymorph of $\mathrm{Co}(\mathrm{acac})_{3}$ has been reported, crystallizing in the monoclinic crystal system with space group $P 2_{1} / c$ (Hon \& Pfluger, 1973; Krüger \& Reynhardt, 1974). We report here a multi-temperature measurement using the temperatures 240,
( $\mathrm{I} a$ ), 210, ( $\mathrm{I} b$ ), 180, ( $\mathrm{I} c), 150,(\mathrm{I} d)$, and 110 K , ( $\mathrm{I} e$ ), which are the same as in our previous publication on $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007).

$\mathrm{Co}(\mathrm{acac})_{3}$ crystallizes as an $\alpha$ polymorph with one independent molecule in the asymmetric unit. No phase transition was observed in the chosen temperature range. The temperature-dependent measurement shows, as expected, a decrease of the anisotropic displacement parameters with decreasing temperature and therefore a decrease of the thermal motion (Fig. 1). The molecule has an approximate noncrystallographic $D_{3}$ symmetry, with r.m.s. deviations between 0.114 [for ( $\mathrm{I} e$ )] and $0.122 \AA$ [for ( $\mathrm{I} a$ )] from ideal symmetry (Pilati \& Forni, 1998).

It is known that it is difficult to reach the precision of fourcircle diffractometer cell parameters from area-detector data (Herbstein, 2000). Two difficulties are noteworthy: (i) the detector is moveable and the detector-crystal distance is highly correlated with the unit-cell volume, and (ii) the precise peak position in rotation direction (time axis) is unknown. In order to overcome some of these difficulties, we carried out temperature-dependent cell-parameter determinations with a fixed detector position and the Phi/Phi-Chi routine (Duisenberg et al., 2000), which is designed to provide accurate peak positions for the reflections. The crystal used for the abovedescribed intensity measurements was different from that used for the crystal structure determinations. The chosen temperature range was from 290 to 110 K , with steps of 20 K and a cooling rate of $120 \mathrm{~K} \mathrm{~h}^{-1}$. The cell axes and the cell volume decrease linearly, the $a$ axis being the most sensitive to change of temperature. The slope ratio $b: c: a$ is approximately 1:1.8:2.7 (Fig. 2). The cell parameters of the known crystal structure determinations of $\mathrm{Co}(\mathrm{acac})_{3}$ at room temperature are summarized in Table 1. It can be clearly seen that the consistency of these data is much worse than the internal consistency given in Fig. 2, which is based on one single crystal and one diffractometer.

The thermal expansion can also be represented by a symmetrical second-rank tensor, $\alpha_{i j}$. Calculations for the tensor components, eigenvalues and eigenvectors were performed using the program $A L P H A$ (Jessen \& Küppers, 1991). 20 reflections were included in the calculation with $\theta$ values calculated from the cell parameters. In a monoclinic unit cell, two off-diagonal terms of the tensor are zero and one
axis is parallel to the crystallographic $b$ axis. The behavior of the thermal expansion is anisotropic, as can be seen in the values for the tensor components of thermal expansion $\alpha_{i j}$
(Table 2). The largest expansion of $\operatorname{Co}(\mathrm{acac})_{3}$ occurs along the crystallographic $a$ axis. This expansion is represented by the eigenvalue of $\alpha_{3}$, which is almost parallel to the crystal-

(1a)

(1b)

(1c)

(1d)

(1e)

Figure 1
Displacement ellipsoid plots and atomic numbering schemes of ( $\mathrm{I} a$ ) at 240 K , ( $\mathrm{I} b$ ) at 210 K , ( $\mathrm{I} c$ ) at 180 K , ( $\mathrm{I} d$ ) at 150 K and ( $\mathrm{I} e$ ) at 110 K . Ellipsoids are drawn at the $50 \%$ probability level. H atoms have been omitted for clarity. All plots are drawn in the same orientation and have the same labeling scheme.

Figure 2

(a) Unit-cell lengths $a$ (middle line), $b$ (lower line) and $c$ (top line), (b) cell angle $\beta$, and ( $c$ ) cell volume as a function of temperature. $(a=13.486+1.32 \times$ $10^{-3} T, b=7.313+0.50 \times 10^{-3} T, c=15.932+0.88 \times 10^{-3} T$ and $V=1552.6+0.357 T$.)
lographic $a$ axis (Table 3). Note that the $\beta$ angle and its change cannot be determined reliably; for example, a relatively large change in $\beta$ from 98.5 to $98.6^{\circ}$ results in an extremely small change in $\theta$ of $0.006^{\circ}$ for the 002 reflection.

The anisotropic thermal expansion tensor might give valuable insight into the strengths of intermolecular interactions (Salud et al., 1998; Kitaigorodsky, 1973), while in other cases such a direct relation has been questioned (Boldyreva et al., 1997). It has been noted that acetylacetonate complexes are linked together in the crystal by pure van der Waals interactions (Alekseev et al., 2006), but in the case of $\alpha-\operatorname{Co}(\mathrm{acac})_{3}$, we detect three different $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ hydrogen bonds of different lengths (Table 4). Methyl H atoms act as hydrogenbond donors and metal-coordinated O atoms as acceptors, with $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ angles between 170 and $180^{\circ}$ for determinations $(\mathrm{I} a)-(\mathrm{I} e)$. The weakest interaction is $\mathrm{C} 11-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ [from $2.78 \AA$ for ( $\mathrm{I} a$ ) to $2.65 \AA$ for (Ie); symmetry code: (iii) $-x,-y+1,-z+1]$, and the vector $\mathrm{C} 11 \cdots \mathrm{O} 1^{\mathrm{iii}}$ forms an angle of $33^{\circ}$ with the crystallographic $a$ axis. The other two interactions, $\mathrm{C} 6-\mathrm{H} 6 B \cdots \mathrm{O}^{\mathrm{ii}}$ and $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$, are shorter [ $2.53 \AA$ for (I $a$ ) to 2.47 for (Ie), and $2.54 \AA$ for ( $\mathrm{I} a$ ) to $2.52 \AA$ for (Ie); symmetry codes: (i) $x, y+1, z$; (ii) $\left.x,-y+\frac{1}{2}, z-\frac{1}{2}\right]$ and form angles of 84 and $72^{\circ}$, respectively, with the crystallographic $a$ axis (Fig. 3). The lengths of these hydrogen bonds


Figure 3
Hydrogen-bond interactions in ( $\mathrm{I} e$ ). The $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ contacts are shown with dashed lines. (a) The $\mathrm{C} 1-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ contact, viewed along the crystallographic $c$ axis, $(b)$ the $\mathrm{C} 6-\mathrm{H} 6 B \cdots \mathrm{O} 6^{\mathrm{ii}}$ contact, viewed along the crystallographic $b$ axis, and (c) the $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$ contact, viewed along the crystallographic $a$ axis. [Symmetry codes: (i) $x, y+1, z$; (ii) $x,-y+\frac{1}{2}$, $z-\frac{1}{2}$; (iii) $-x,-y+1,-z+1$.]
are linearly dependent on the temperature (Fig. 4). Interestingly, the longest distance, $\mathrm{C} 11 \cdots \mathrm{O} 1^{\mathrm{iii}}$, has the largest slope. This hydrogen bond has the smallest angle to the crystallographic $a$ axis, which has also the largest temperature


Figure 4
$D \cdots A$ distances for the hydrogen-bond contacts $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$ (middle line), $\mathrm{C} 6-\mathrm{H} 6 B \cdots \mathrm{O} 6^{\text {ii }}$ (lower line) and $\mathrm{C} 11-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ (top line) as a function of temperature $\left[\mathrm{C} 1 \cdots \mathrm{O} 4^{\mathrm{i}}=3.482+0.13 \times 10^{-3} T, \mathrm{C} 6 \cdots \mathrm{O}^{6 \mathrm{ii}}=\right.$ $3.399+0.38 \times 10^{-3} T$ and $\mathrm{C} 11 \cdots \mathrm{O} 1^{\mathrm{iii}}=3.530+0.89 \times 10^{-3} T$; symmetry codes: (i) $x, y+1, z$; (ii) $x,-y+\frac{1}{2}, z-\frac{1}{2}$; (iii) $-x,-y+1,-z+1$.]


Figure 5
Eigenvalues of translation $T_{i}$ (top) and libration $L_{i}$ (bottom) tensors from rigid-body analyses (PLATON; Spek, 2003) as a function of temperature. $\left(T_{1}=1.68 \times 10^{-3}+0.15 \times 10^{-3} T, T_{2}=0.95 \times 10^{-3}+0.13 \times 10^{-3} T, T_{3}=\right.$ $0.76 \times 10^{-3}+0.11 \times 10^{-3} T, L_{1}=-0.67+0.06 T, L_{2}=-0.86+0.02 T$ and $\left.L_{3}=-0.22+0.02 T.\right)$

Figure 6


Peanut plots of $(\mathrm{I} a)$ at 240 K , ( $\mathrm{I} b$ ) at 210 K , ( $\mathrm{I} c)$ at 180 K , ( $\mathrm{I} d$ ) at 150 K and ( $\mathrm{I} e$ ) at 110 K , showing the difference between the measured displacement parameters and the parameters obtained by rigid-body analyses using the program THMA11 (Schomaker \& Trueblood, 1998). A scale factor of 3.08 was used for the r.m.s. surfaces. Darker lines (blue in the electronic version) indicate positive differences and lighter lines (purple in the electronic version) negative differences. All plots are drawn in the same orientation and have the same labelling scheme.
dependence of the cell parameters (see above). The other two hydrogen bonds are approximately in the $b c$ plane and have a smaller temperature dependence.

The eigenvalues of the translation $\mathbf{T}$ and libration $\mathbf{L}$ tensors obtained from rigid-body analyses (PLATON; Spek, 2003) show that the decrease of the thermal motion is linear in the whole temperature range (Fig. 5). The agreement factors of the rigid-body analyses are rather high $\{R=0.104-0.129 ; R=$ $\left.\left[\sum\left(U_{\text {obs }}-U_{\text {calc }}\right)^{2} / \sum U_{\text {obs }}^{2}\right]^{1 / 2}\right\}$. This situation is similar to that for $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$, where the agreement factors $(R=0.155-$ 0.167 ) are even higher. The corresponding internal motion is visualized by difference plots (Hummel et al., 1990) between the observed displacement parameters and the rigid-body models (Fig. 6). The nonrigid behavior is mainly expressed by only one acac ligand (C11-C15). This behavior is also observed in the aluminium compound.

Considering the displacement parameters at a particular temperature, the internal motions are also reflected in a relatively large variation of the bond lengths. For example, the $\mathrm{Co}-\mathrm{O}$ distances range between 1.8746 (18) and 1.8892 (17) $\AA$ at 240 K , viz. determination (Ia). The corresponding Al-O distances in $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$ for 240 K are 1.8712 (13)1.8862 (12) Å and show a similar variation (von Chrzanowski et al., 2007). Because octahedral, low-spin $\mathrm{Co}^{\text {III }}$ with electron configuration $d^{6}$ does not express Jahn-Teller distortions (Wiberg, 1985), this variation can only be explained by internal thermal motion. The thermal motion also contributes to a shortening of the $\mathrm{Co}-\mathrm{O}$ distances (Table 5) of the crystal structure determinations compared with the value of $1.899 \AA$ obtained for $\mathrm{Co}(\mathrm{acac})_{3}$ from high-level density functional theory ( $6-31 \mathrm{G}^{*}$ for $\mathrm{C}, \mathrm{H}, \mathrm{O}$ and triple- $\zeta$ for Co ) calculations (Diaz-Acosta et al., 2003).

## Experimental

Dark-green crystals were obtained by slow evaporation of a solution of the commercially available material (Aldrich) in ethyl acetate at room temperature.

## Compound (Ia), at 240 K

Crystal data
$\left[\mathrm{Co}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$
$M_{r}=356.25$
Monoclinic, $P 2_{1} / c$
$a=13.8094$ (7) $\AA$
$b=7.4331$ (4) $\AA$
$c=16.1484(9) \AA$
$\beta=98.430(3)^{\circ}$

## Data collection

Nonius KappaCCD diffractometer Absorption correction: multi-scan (SADABS; Sheldrick, 2002) $T_{\text {min }}=0.47, T_{\text {max }}=0.95$

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.040$
$w R\left(F^{2}\right)=0.091$
$S=1.04$
3762 reflections

## Compound (Ib), at 210 K

## Crystal data

$\left[\mathrm{Co}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$
$M_{r}=356.25$
Monoclinic, $P 2_{1} / c$
$a=13.7708$ (8) $\AA$
$b=7.4195$ (3) $\AA$
$c=16.1242(7) \AA$
$\beta=98.455(2)^{\circ}$

$$
\begin{aligned}
& V=1639.66(15) \AA^{3} \\
& Z=4 \\
& \text { Mo } K \alpha \text { radiation } \\
& \mu=1.07 \mathrm{~mm}^{-1} \\
& T=240(2) \mathrm{K} \\
& 0.23 \times 0.15 \times 0.05 \mathrm{~mm}
\end{aligned}
$$

18651 measured reflections 3762 independent reflections 2578 reflections with $I>2 \sigma(I)$ $R_{\text {int }}=0.044$

## 205 parameters

H -atom parameters constrained
$\Delta \rho_{\text {max }}=0.28 \mathrm{e}_{\AA^{-3}}$
$\Delta \rho_{\text {min }}=-0.33 \mathrm{e}^{-3}$

$$
\begin{aligned}
& V=1629.54(14) \AA^{3} \\
& Z=4 \\
& \text { Mo } K \alpha \text { radiation } \\
& \mu=1.08 \mathrm{~mm}^{-1} \\
& T=210(2) \mathrm{K} \\
& 0.23 \times 0.15 \times 0.05 \mathrm{~mm}
\end{aligned}
$$

## Data collection

Nonius KappaCCD diffractometer Absorption correction: multi-scan (SADABS; Sheldrick, 2002) $T_{\text {min }}=0.52, T_{\text {max }}=0.95$

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.039$
$w R\left(F^{2}\right)=0.087$
$S=1.07$
3741 reflections

## Compound (Ic), at 180 K

## Crystal data

$\left[\mathrm{Co}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$
$M_{r}=356.25$
Monoclinic, $P 2_{1} / c$
$a=13.7338$ (6) $\AA$
$b=7.4070(3) \AA$
$c=16.0959$ (7) $\AA$
$\beta=98.467$ (1) ${ }^{\circ}$

## Data collection

Nonius KappaCCD diffractometer Absorption correction: multi-scan (SADABS; Sheldrick, 2002)
$T_{\text {min }}=0.51, T_{\text {max }}=0.95$

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.036$
$w R\left(F^{2}\right)=0.079$
$S=1.04$
3713 reflections

## Compound (Id), at 150 K

## Crystal data

$\left[\mathrm{Co}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$
$M_{r}=356.25$
Monoclinic, $P 2_{1} / c$
$a=13.6927$ (6) $\AA$
$b=7.3920$ (3) $\AA$
$c=16.0641$ (6) A
$\beta=98.499(2)^{\circ}$
Data collection
Nonius KappaCCD diffractometer Absorption correction: multi-scan (SADABS; Sheldrick, 2002)
$T_{\text {min }}=0.47, T_{\text {max }}=0.95$

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.035$
$w R\left(F^{2}\right)=0.080$
$S=1.03$
3698 reflections

## Compound (Ie), at 110 K

Crystal data
$\left[\mathrm{Co}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$
$M_{r}=356.25$
Monoclinic, $P 2_{1} / c$
$a=13.6376$ (5) $\AA$
$b=7.3758(3) \AA$
$c=16.0446(9) \AA$
$\beta=98.594(2)^{\circ}$

18555 measured reflections 3741 independent reflections 2680 reflections with $I>2 \sigma(I)$ $R_{\text {int }}=0.042$

## 205 parameters

H -atom parameters constrained
$\Delta \rho_{\text {max }}=0.28 \mathrm{e}^{-3}$
$\Delta \rho_{\min }=-0.35 \mathrm{e} \mathrm{A}^{-3}$
$V=1619.53(12) \AA^{3}$
$Z=4$
Mo $K \alpha$ radiation
$\mu=1.08 \mathrm{~mm}^{-1}$
$T=180$ (2) K
$0.23 \times 0.15 \times 0.05 \mathrm{~mm}$

18433 measured reflections 3713 independent reflections 2765 reflections with $I>2 \sigma(I)$ $R_{\text {int }}=0.039$

205 parameters
H -atom parameters constrained
$\Delta \rho_{\text {max }}=0.29 \mathrm{e}_{\AA^{-3}}$
$\Delta \rho_{\text {min }}=-0.35 \mathrm{e}^{-3}$
$Z=4$
Mo $K \alpha$ radiation
$\mu=1.09 \mathrm{~mm}^{-1}$
$T=150(2) \mathrm{K}$
$0.23 \times 0.15 \times 0.05 \mathrm{~mm}$

18370 measured reflections
3698 independent reflections
2812 reflections with $I>2 \sigma(I)$ $R_{\text {int }}=0.043$

205 parameters
H-atom parameters constrained
$\Delta \rho_{\text {max }}=0.31 \mathrm{e}_{\AA}^{-3}$
$\Delta \rho_{\min }=-0.42 \mathrm{e}^{-3}$
$V=1595.78(12) \AA^{3}$
$Z=4$
Mo $K \alpha$ radiation
$\mu=1.10 \mathrm{~mm}^{-1}$
$T=110$ (2) K
$0.23 \times 0.15 \times 0.05 \mathrm{~mm}$

## Data collection

Nonius KappaCCD diffractometer Absorption correction: multi-scan (SADABS; Sheldrick, 2002)
$T_{\text {min }}=0.57, T_{\text {max }}=0.95$

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.029$
$w R\left(F^{2}\right)=0.065$
$S=1.06$
3656 reflections

205 parameters
27696 measured reflections 3656 independent reflections 3018 reflections with $I>2 \sigma(I)$ $R_{\text {int }}=0.040$

H -atom parameters constrained
$\Delta \rho_{\text {max }}=0.32 \mathrm{e} \AA^{-3}$
$\Delta \rho_{\text {min }}=-0.37 \mathrm{e}^{-3}$

Table 1
Cell parameters $\left(\AA,{ }^{\circ}\right)$ of the known room-temperature $\mathrm{Co}(\mathrm{acac})_{3}$ crystal structure determinations.

| $a$ | $b$ | $c$ | $\beta$ | Reference |
| :--- | :--- | :--- | :--- | :--- |
| 14.160 | 7.480 | 16.430 | 98.68 | Padmanabhan (1958) |
| $14.25(3)$ | $7.50(2)$ | $16.38(3)$ | $99(1)$ | Shkol'nikova (1959) |
| $13.951(9)$ | $7.470(5)$ | $16.222(11)$ | $98.48(8)$ | Hon \& Pfluger (1973) |
| $13.90(7)$ | $7.47(4)$ | $16.21(7)$ | $98.4(1)$ | Krüger \& Reynhardt (1974) |
| $13.8663(15)$ | $7.4599(5)$ | $16.1874(9)$ | $98.409(5)$ | Present study |

Table 2
Tensor components of the thermal expansion $\left(10^{-6} \mathrm{~K}^{-1}\right)$.
Standard uncertainties are underestimated because calculated $\theta$ values were used (see Comment).

| $T(\mathrm{~K})$ | $\alpha_{11}$ | $\alpha_{22}$ | $\alpha_{33}$ | $\alpha_{12}$ | $\alpha_{13}$ | $\alpha_{23}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $290-270$ | $80.79(3)$ | $76.54(4)$ | $53.01(17)$ | 0 | $-4.34(12)$ | 0 |
| $270-250$ | $94.35(2)$ | $75.41(3)$ | $57.17(13)$ | 0 | $-3.63(9)$ | 0 |
| $250-230$ | $94.51(3)$ | $75.45(4)$ | $53.85(17)$ | 0 | $-4.30(11)$ | 0 |
| $230-210$ | $101.56(2)$ | $72.40(3)$ | $59.04(13)$ | 0 | $5.25(9)$ | 0 |
| $210-190$ | $89.56(3)$ | $67.08(4)$ | $49.23(18)$ | 0 | $2.46(12)$ | 0 |
| $190-170$ | $107.55(2)$ | $52.85(3)$ | $56.87(14)$ | 0 | $10.89(10)$ | 0 |
| $170-150$ | $100.83(2)$ | $74.01(3)$ | $57.90(12)$ | 0 | $5.32(8)$ | 0 |
| $150-130$ | $108.88(2)$ | $50.06(3)$ | $52.58(10)$ | 0 | $9.17(7)$ | 0 |
| $130-110$ | $96.35(2)$ | $61.16(3)$ | $47.50(11)$ | 0 | $7.34(7)$ | 0 |

Table 3
Eigenvalues $\left(10^{-6} \mathrm{~K}^{-1}\right)$ and eigenvectors of the thermal expansion tensor.

| $T(\mathrm{~K})$ | Principal axis | Eigenvalue | Eigenvector |  |  |
| :--- | :--- | :---: | :---: | :--- | :--- |
| $290-270$ | $\alpha_{1}$ | 52.3 | 0.1513 | 0 | 0.9885 |
|  | $\alpha_{2}$ | 76.5 | 0 | 1.0 | 0 |
|  | $\alpha_{3}$ | 81.5 | -0.9886 | 0.0 | 0.1508 |
| $270-250$ | $\alpha_{1}$ | 56.8 | 0.0965 | 0.0 | 0.9953 |
|  | $\alpha_{2}$ | 75.4 | 0 | 1.0 | 0 |
|  | $\alpha_{3}$ | 94.7 | -0.9954 | 0 | 0.0960 |
| $250-230$ | $\alpha_{1}$ | 53.4 | 0.1043 | 0 | 0.9945 |
|  | $\alpha_{2}$ | 75.5 | 0 | 1.0 | 0 |
|  | $\alpha_{3}$ | 95.0 | -0.9946 | 0 | 0.1039 |
| $230-210$ | $\alpha_{1}$ | 58.4 | -0.1211 | 0 | 0.9926 |
|  | $\alpha_{2}$ | 72.4 | 0 | -1.0 | 0 |
|  | $\alpha_{3}$ | 102.2 | 0.9927 | 0 | 0.1207 |
| $210-190$ | $\alpha_{1}$ | 49.1 | -0.0610 | 0 | 0.9981 |
|  | $\alpha_{2}$ | 67.1 | 0 | -1.0 | 0 |
|  | $\alpha_{3}$ | 89.7 | 0.9982 | 0 | 0.0606 |
| $190-170$ | $\alpha_{1}$ | 52.8 | 0 | -1.0 | 0 |
|  | $\alpha_{2}$ | 54.6 | -0.2019 | 0 | 0.9794 |
|  | $\alpha_{3}$ | 109.8 | 0.9795 | 0 | 0.2015 |
| $170-150$ | $\alpha_{1}$ | 57.3 | -0.1215 | 0 | 0.9926 |
|  | $\alpha_{2}$ | 74.0 | 0 | -1.0 | 0 |
| $150-130$ | $\alpha_{3}$ | $\alpha_{1}$ | 101.5 | 0.9926 | 0 |
|  | $\alpha_{2}$ | 50.1 | 0 | 0.1211 |  |
|  | $\alpha_{3}$ | 51.1 | -0.1570 | 0 | 0 |
| $130-110$ | $\alpha_{1}$ | 110.3 | 0.9876 | 0 | 0.9876 |
|  | $\alpha_{2}$ | 46.4 | -0.1456 | 0 | 0.9897 |
|  | $\alpha_{3}$ | 61.2 | 0 | -1.0 | 0 |
|  |  | 97.4 | 0.9894 | 0 | 0.1453 |

Table 4
Hydrogen-bond geometry $\left({ }^{\circ},^{\circ}\right)$ for $(\mathrm{I} a)(240 \mathrm{~K})$ to $(\mathrm{I} e)(110 \mathrm{~K})$.

|  | $D-\mathrm{H} \cdots A$ | $D-\mathrm{H}$ | $\mathrm{H} \cdots A$ | $D \cdots A$ | $D-\mathrm{H} \cdots A$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| (Ia) | $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$ | 0.97 | 2.54 | 3.513 (3) | 178 |
| (Ib) | $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$ | 0.97 | 2.54 | 3.512 (3) | 179 |
| (Ic) | $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4{ }^{\mathrm{i}}$ | 0.98 | 2.53 | 3.507 (3) | 178 |
| (Id) | $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$ | 0.98 | 2.52 | 3.501 (3) | 179 |
| (Ie) | $\mathrm{C} 1-\mathrm{H} 1 A \cdots \mathrm{O} 4^{\mathrm{i}}$ | 0.98 | 2.52 | 3.497 (2) | 180 |
| (Ia) | C6-H6B . $\mathrm{O}^{\text {ii }}$ | 0.97 | 2.53 | 3.492 (3) | 172 |
| (Ib) | C6-H6B..O6 ${ }^{\text {ii }}$ | 0.97 | 2.52 | 3.478 (3) | 172 |
| (Ic) | C6-H6B . $\mathrm{O}^{\text {ii }}{ }^{\text {ii }}$ | 0.98 | 2.50 | 3.467 (3) | 172 |
| (Id) | C6- $\mathrm{H} 6 \mathrm{~B} \cdots \mathrm{O}^{\text {ii }}$ | 0.98 | 2.48 | 3.455 (3) | 171 |
| (Ie) | C6-H6B.. $\mathrm{O}^{\text {ii }}$ | 0.98 | 2.47 | 3.442 (2) | 171 |
| (Ia) | $\mathrm{C} 11-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ | 0.97 | 2.78 | 3.743 (4) | 175 |
| (Ib) | $\mathrm{C} 11-\mathrm{H} 118 \cdots \mathrm{O} 1^{\text {iii }}$ | 0.97 | 2.75 | 3.716 (4) | 175 |
| (Ic) | $\mathrm{C} 11-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ | 0.98 | 2.71 | 3.690 (3) | 174 |
| (Id) | $\mathrm{C} 11-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ | 0.98 | 2.68 | 3.660 (3) | 177 |
| (Ie) | $\mathrm{C} 11-\mathrm{H} 11 B \cdots \mathrm{O} 1^{\text {iii }}$ | 0.98 | 2.65 | 3.629 (2) | 178 |

Symmetry codes: (i) $x, y+1, z$; (ii) $x,-y+\frac{1}{2}, z-\frac{1}{2}$; (iii) $-x,-y+1,-z+1$.

Table 5
Co-O bond lengths ( $\AA$ ) for ( $\mathrm{I} a$ ) ( 240 K ) to ( $\mathrm{I} e$ ) ( 110 K ).

| Co-O | (I $a)$ | $(\mathrm{I} b)$ | $(\mathrm{I} c)$ | $(\mathrm{I} d)$ | $(\mathrm{I} e)$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| Co1-O5 | $1.8746(18)$ | $1.8748(17)$ | $1.8762(16)$ | $1.8766(15)$ | $1.8770(12)$ |
| Co1-O6 | $1.8798(17)$ | $1.8815(16)$ | $1.8813(15)$ | $1.8801(14)$ | $1.8827(12)$ |
| Co1-O3 | $1.8826(17)$ | $1.8834(16)$ | $1.8851(15)$ | $1.8835(14)$ | $1.8859(12)$ |
| Co1-O4 | $1.8841(17)$ | $1.8846(16)$ | $1.8853(15)$ | $1.8856(14)$ | $1.8865(12)$ |
| Co1-O2 | $1.8872(17)$ | $1.8878(16)$ | $1.8893(15)$ | $1.8899(14)$ | $1.8906(12)$ |
| Co1-O1 | $1.8892(17)$ | $1.8884(16)$ | $1.8902(15)$ | $1.8912(14)$ | $1.8913(12)$ |

The crystal structure determinations of (Ia)-(Ie) were all carried out on the same crystal, which was mounted on a glass capillary with perfluoro polyether oil (Kottke \& Stalke, 1993). The X-ray intensities of (Ia)-(Id) were obtained with an exposure time of 40 s per frame and a rotation angle of $1^{\circ} .258 \varphi$ and $125 \omega$ scans were measured. The X-ray intensities of (Ie) were obtained with two different exposure times and rotation angles of $1^{\circ} .258 \varphi$ and $323 \omega$ scans were measured with an exposure time of 60 s per frame and $173 \varphi$ scans with an exposure time of 12 s per frame.

For the structure solution of ( $\mathrm{I} a$ ), coordinates were taken from those of $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007); subsequently, coordinates were taken from ( $\mathrm{I} a$ ) for ( $\mathrm{I} b$ ), from ( $\mathrm{I} b$ ) for ( $\mathrm{I} c$ ), from (I $c$ ) for ( $\mathrm{I} d$ ) and from ( $\mathrm{I} d$ ) for ( $\mathrm{I} e$ ). All H atoms were introduced in geometrically idealized positions, refined with a riding model and subsequently confirmed in a difference Fourier map. Their $U_{\text {iso }}(\mathrm{H})$ values were set at $1.2 U_{\text {eq }}(\mathrm{C})$ for H atoms of the central CH groups and at $1.5 U_{\text {eq }}(\mathrm{C})$ for methyl H atoms.

Because the above-mentioned oil mounting technique is not suitable for ambient temperatures, a second, different, crystal was selected and mounted with Super Glue for the cell-parameter
determinations. The detector position was kept fixed and the Phi/PhiChi routine (Duisenberg et al., 2000) was used.

For all determinations, data collection: COLLECT (Nonius, 1999); cell refinement: PEAKREF (Schreurs, 2005); data reduction: EVALCCD (Duisenberg et al., 2003) and SADABS (Sheldrick, 2002); program(s) used to refine structure: SHELXL97 (Sheldrick, 1997); molecular graphics: PLATON (Spek, 2003); software used to prepare material for publication: manual editing of SHELXL97 output.

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Supplementary data for this paper are available from the IUCr electronic archives (Reference: GZ3091). Services for accessing these data are described at the back of the journal.

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